$F_{t_y}$ 



Using the Ramberg-Osgood (Hill) stress-strain formulation

0.002

$$
\varepsilon_{u}^{\prime} = \varepsilon_{u} - \frac{F_{tu}}{E} = 0
$$
  

$$
n = \log_{e} \left( \frac{\varepsilon_{u}^{\prime}}{a \cdot \omega^{2}} \right) / \log_{e} \left( \frac{F_{tu}}{F_{t_{y}}} \right) - 0
$$
  

$$
\varepsilon_{m} = \frac{F_{m}}{E} + \varepsilon_{u}^{\prime} \left( \frac{F_{m}}{F_{tu}} \right)^{n} - 0
$$
  
is a result of section (b x h) the distance

For a rectangular ser<br>bending moment is:-88

$$
M_{C_{e_2}} = \frac{b h^2}{12} (2 f_m + f_o) \qquad -\textcircled{4}
$$

By integration of 3 over the section, yields the<br>plastic bending moment:-

$$
M_{max} = \frac{bh^2 f_m}{g} \left\{ \frac{1}{3} \left( \frac{f_m}{E} \right)^2 + \varepsilon_u' \left( \frac{n+1}{n+2} \right) \frac{f_m}{f_{\text{cut}}} \right)^n / \frac{f_m}{E} + \left( \varepsilon_u' \right)^2 \left( \frac{n}{2n+1} \right) \left( \frac{f_m}{F_{\text{cut}}} \right)^{2n} \left( \frac{f_m}{E} \right) + \left( \varepsilon_u' \right)^2 \left( \frac{n}{2n+1} \right) \left( \frac{f_m}{F_{\text{cut}}} \right)^{2n} \left( \frac{f_m}{E} \right) \n\frac{f_m}{F_m} = \left( \frac{6}{\varepsilon_n^2} \right) \cdot \left( \frac{1}{3} \left( \frac{f_m}{E} \right)^2 + \varepsilon_u' \left( \frac{n+1}{n+2} \right) \left( \frac{f_m}{F_{\text{cut}}} \right)^n \left( \frac{f_m}{E} \right) + \left( \varepsilon_u' \right)^2 \cdot \left( \frac{n}{\varepsilon_{n+1}} \right) \cdot \left( \frac{f_m}{F_{\text{cut}}} \right)^{2n} \right) - 2
$$

Thus from the material parameters  
\n• 
$$
F_{t_u}
$$
  
\n•  $E$   
\n•  $E_u$ 

the Coggone stresses fo and  $f_m$  can be determined. aesommen.<br>The key is the sentence from Peery: -<br>"Cozzone has shown that it is sufficiently<br>accurate to calculate fo for a rectangular cross section and to use this value for all cross sections

$$
\frac{3}{\cos}
$$

$$
E_t = \left[\frac{1}{\sqrt{\frac{1}{E} + \left(\frac{n}{F_{tu}}\right)} \cdot \varepsilon_u' \cdot \left(\frac{f}{F_{tu}}\right)^{n-1}}\right] \quad 2 \, \odot
$$

Secant Modulus: -

$$
E_{s} = \frac{1}{\left[\frac{1}{E} + \left(\frac{\varepsilon_{u}}{F_{t_{u}}}\right) \cdot \left(\frac{f}{F_{t_{u}}}\right)^{n-1}\right]} \geq 0
$$

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$ 

For compression buckling and shear buckling<br>in the plastic range, the buckling formula is<br>written as follows:-

$$
f_{cr} = k \cdot \gamma_p \cdot \mathcal{E} \cdot \left(\frac{t}{b}\right)^2
$$

When:  
\nFor compression buckling, 
$$
\gamma_p = \gamma_c = \left(\frac{E_t}{E}\right)^{1/2}
$$
  
\nand,  
\nFor shear buckling,  $\gamma_p = \gamma_s = \left(\frac{G_t}{G}\right)^{1/2}$   
\nwhere :-  
\n $G = \frac{E}{2(1+\nu)}$ 

$$
k = k_c
$$
 or  $k_s \cdot \frac{\pi^2}{12(1-\nu^2)}$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

$$
\frac{4}{\cos 2}
$$

Alternate stress-strain curve formulation.	
Using:	$E_g' = E_g - \frac{E_g}{E} = 0.002$ (by def.)
$n = same as E_gu$	
then:	$E = \frac{F}{E} + E_g'(\frac{F}{E_g})^n$
Note:	It is common in compression analysis to substitute E_g for E_g.
Thus we can write the 7c and 7s terms as follows:	
$0c = \left[ \frac{1}{1 + (\frac{0.002 E_n}{E_g})(\frac{f_{cr}}{E_g})^{n-1}} \right]^{1/2}$	
$0c = \left[ \frac{1}{1 + (\frac{0.002 E_n}{E_g})(\frac{f_{cr}}{E_g})^{n-1}} \right]^{1/2}$	
$0s = \left[ \frac{1}{1 + (\frac{0.002 G_n}{E_g}) (\frac{f_{cr}}{E_g})^{n-1}} \right]^{1/2}$	

where 
$$
F_{s_3} \doteq 0.55 F_{c_3}
$$

 $5\overline{)}$  $coz$ 

Much data exists for compression failure in terms<br>of the f<sub>ay</sub> stress where the 70% secant madulus<br>cuts the stress-strain curve.

Using Egu. 
$$
\odot
$$
 and rearsanging, we can write  

$$
f_{0,7} = F_{cy} \left[ 214.286 \cdot \frac{F_{cy}}{E_c} \right]^{n-1}
$$

ĥ  $coz$ 



From the Bernoulli-Euler theory of bending, also known as the Engineers' Theory of Bending (ETB) in which the assumption is made that plane sections remain plane, the strain distribution is linear. V

The maximum strain, at the outer fibres of the beam, is established by the maximum allowable compressive stress. The compressive stress is limited by the instability or collapse mode of the compression member located at the extreme fibre of the beam. This maximum compressive stress, and associated strain, is denoted as  $f_m$  and  $\epsilon_m$  respectively.

Thus at a distance y from the neutral axis (NA)<br> $E(y) = \frac{\epsilon_m}{h/z} \cdot y$ 

W This is apparantly also referred to as the Bernoulli-<br>Hypothesis, in some academic circles. (Ref. RM Rivello)

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ 

 $\overline{7}$  $coz$ 

 $\alpha$ 

$$
y = \frac{h_{/z}}{\varepsilon_m} \cdot \varepsilon(y)
$$

$$
= \left(\frac{h}{\varepsilon_m}\right) \cdot \left[\frac{f}{\varepsilon} + \varepsilon'_g \left(\frac{f}{F_{cg}}\right)^n\right] \geq 0
$$

Differentiating w.r.t. 
$$
f
$$
, yields:-  
 
$$
dy = \left(\frac{h}{eE_m}\right) \left[\frac{1}{E} + \frac{nE_y'}{E_y'} f^{n-1}\right] dP
$$

Determining the moment carried by  $\frac{1}{2}$  of the section, yields : $h/a$ 

$$
H_{R}M_{m} = b \int_{0}^{m_{R}} f \cdot g \cdot dy
$$

Substituting 
$$
Eg
$$
 and  $\Theta$  and  $\Theta$  and change the integration bounds to suit, yields :-  
\n
$$
\frac{1}{2}M_m = b \left(\frac{h}{2E_m}\right)^2 \left[\frac{f}{E} + E_g'(\frac{f}{E_g})^n\right] \cdot \left[\frac{f}{E} + n \frac{E_g' (f)}{E_g}\right]^n d\theta
$$

$$
=b\left(\frac{h}{2\epsilon_m}\right)\left[\frac{\int_{\epsilon_m}^{\epsilon_m}f(x)\,dx}{\int_{\epsilon_m}^{\epsilon_m}f(x)\,dx}\right]^{2}+\left(n+1\right)\epsilon_{s}\left(\frac{f}{E}\right)\left(\frac{f}{E_{s}}\right)^{n}\right]d\ell
$$

Integrating w.r.t. "f", gields:-

$$
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$$

$$
H_{m} = b \left( \frac{h}{2\epsilon_{m}} \right)^{2} \left[ \frac{1}{3} \frac{f^{3}}{\epsilon^{2}} + \frac{n \epsilon_{y}^{'2}}{(2n+1)\epsilon_{cy}^{2n}} f^{2n+1} + \epsilon_{y}^{'n+1} \frac{f^{n+2}}{n+2} f^{2n} \right]_{0}^{2n}
$$
  
\nSubstituting the integration bounds into the function  
\nyields:-  
\n
$$
M_{m} = \frac{bh^{2}}{\epsilon} \frac{F_{m}}{\epsilon_{m}^{2}} \left\{ \frac{f_{m}}{\epsilon} \right\}^{2} + \left( \epsilon_{y}^{'2} \frac{n}{(2n+1)} \left( \frac{f_{m}}{f_{cy}} \right)^{2n} + \epsilon_{y}^{'n} \left( \frac{n+1}{f_{cy}} \right)^{n} \left( \frac{f_{m}}{f_{cy}} \right)^{n} \right\}
$$
  
\nwhich is identical to Eqn. 6, but using the function terms  $\epsilon_{y}^{'}$  and  $F_{cy}$ .

In accordance with the derivation on page **1-COZ**, the  $\sigma_{0.7}$  can also be written as follows in terms of the ultimate stress.

$$
\sigma_{0.7} = [(0.4286 \cdot \sigma_{\text{tu}}^{n})/(E_{\text{c}} \cdot \varepsilon_{\text{u}})]^{(1/(n-1))}
$$

## **Additional Observations**

On inserted page 4-COZ above, a statement is made that the material shape power coefficient "n" can be calculated as per the original "full" definition of the stress-strain curve defined on page 1-COZ. From further reading of an aircraft manual from a firm that shall remain nameless, but for whom Prof Niu was a researcher, it has become evident that they pay particular attention to the "knee" of the  $\sigma - \varepsilon$  curve and in fact define the inelastic portion of the curve by means of two zones. The first zone extends from their definition of the proportional limit, 0.0001 to 0.002 permanent set and the second zone extends from 0.002 to **′u**. The usual definition of the proportional limit is 0.001 permanent set strain. It is over the first inelastic zone that the "n" for the formula given on page 4-COZ is intended. Using the overall or "full" definition of "n" does provide the correct inelastic correction outside the permanent set strain range of 0.0001 to 0.002.

This refined usage of the σ-ε curve is extended into the manner in which they calculate the material Cozzone stresses. They also correct for sections that are not rectangular, as asserted by the original Cozzone assumptions.

As a caveat to what the unnamed Big Aircraft Company is doing, it must be noted that many, many other very successful aircraft companies use the method as originally postulated by Cozzone (and presented herein), and is considered to be valid and accurate. This author has also used the method to successfully design and analyse bending components that have been tested to destruction. The method has also been used by the author to "calibrate" aluminium beam section work published on the Cornell University (USA) website. This work was done by the Civil Engineering (Built Environment) Faculty. Their theoretical analysis was performed using inelastic FE Analysis and calibrated against collapse failure laboratory tests. There are additional notes elsewhere in these notes on the

## analysis of I-, C-, and Z-section flanged beams

Another fact that comes to mind here is the question of material strength properties and allowables. The material strength data used by companies, and in the Mil-Handbooks and more recently the MMPDS volumes, are statistically guaranteed values for the alloy and/or from the approved company supplier. When **laboratory tests** are being done, specific **specimen material properties are required** if **accurate** correlation between analysis and test is required.

When only the  $f_0$  Cozzone stress is required at rupture, then a slightly reduced version of the formula on p 2-COZ can be written as follows: -

$$
f_{\text{ou}}/\sigma_{\text{tu}} = (6/\epsilon_{\text{u}}^{2}) \cdot (1/3 \cdot (\sigma_{\text{tu}}/E)^{2} + \epsilon_{\text{u}}' \cdot ((n+1)/(n+2)) \cdot (\sigma_{\text{tu}}/E) + (n/(2 \cdot n+1)) \cdot (\epsilon_{\text{u}}')^{2} - 2
$$

## **10.1.1 Inelastic Buckling Correction**

Up to now inelastic buckling calculations were determined in most American textbooks with the use of dimensionless  $F_{cr}/F_{0.7}$  curves. The introduction of the spreadsheet as an engineering tool has made it possible to directly use the 4 key strength parameters of a material to solve for inelastic buckling stress.

The Hill stress-strain formulation from the 4 parameters  $\sigma_{tu}$ ,  $\sigma_{tv}$ ,  $\varepsilon_{u}$ , **E**<sub>c</sub> yields the material exponent "**n**" and the ultimate material "set" value of **'u** .

In general, inelastic buckling formulas of all types have the following form: -

$$
\sigma_{cr}/\eta_c = \mathbf{k}_c \cdot \mathbf{E}_c \cdot (t/b)^2 \cdots \cdots \cdots \cdots \cdots (A)
$$

For the given geometry of the panel the RHS of (A) is a constant, which we will call **D2**. So: -

$$
D_2 = k_c \cdot E_c \cdot (t/b)^2
$$

The  $\eta_c$  term in turn is a function of the material inelastic parameters and  $\sigma_{cr}$  itself, so an iterative process is required to calculate  $\sigma_{cr}$  that matches the buckling equation (A).

In many of the classical textbooks on airframe analysis the value of  $\eta_c$  is given as: -

**<sup>c</sup> = Et/E<sup>c</sup> ------------------------------(B)** or, **<sup>c</sup> = Es/E<sup>c</sup> ------------------------------(B′)**

While in others it is given as: -

$$
\eta_c
$$
 =  $[E_t/E_c]^{\frac{1}{2}}$ -----  
-----(C)

And more confusingly, some company manuals have fine-tuned these correction functions to even more complex functions based on the stress-strain curve.

For the purposes of these notes and for illustration only, the inelastic correction (B) will be used: -

Thus: - 
$$
E_t/E_c
$$
 = {1/[1+( $\varepsilon'_{u}$ - $E_c$ - $n/\sigma_{cu}$ )( $\sigma_{cr}/\sigma_{cu}$ )<sup>n-1</sup>]}

Substituting into Equ. (A) yields: -

$$
\sigma_{cr} \cdot [1 + (\varepsilon'_{u} \cdot E_{c} \cdot n/\sigma_{cu}) \cdot (\sigma_{cr}/\sigma_{cu})^{n-1}] = D_2
$$

Thus: -

$$
\sigma_{\rm cr} + [(\varepsilon_{\rm u} E_{\rm c} \cdot n/\sigma_{\rm cu}) \cdot (1/\sigma_{\rm cu})^{n-1}] \cdot \sigma_{\rm cr}^{n} = D_2
$$

The terms in the  $\lceil \cdot \rceil$  brackets make up a constant that will be called  $\mathbf{D}_1$ , as follows: -